

# What's the entropy of gravity?

Francesca Vidotto

Università di Pavia and Centre de Physique Théorique, Marseille

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## things you know

Yang-Mills Theory  
electric field, vector potential

angular momentum  
 $L^2|j, m\rangle = \hbar^2 j(j+1)|j, m\rangle$

states

Hilbert space

## Loop Quantum Gravity

Gravity as a  $SU(2)$  gauge theory  
 $(E^{ai}(x), A_a^i(x))$

$$q^{ab}(X) = \frac{E^{ai}(X)E^{bj}(X)}{q(X)}$$

area

$$E_{ie}E_e^i |s\rangle = (8\pi\gamma G\hbar)^2 j_e(j_e + 1) |s\rangle$$

$$\sqrt{q(x)} |s\rangle = \sum_{n \in N(s)} \nu_n \delta(x, x_n) |s\rangle$$

$$|s\rangle := |\Gamma, j_e, \nu_n\rangle$$

$$\mathcal{H}_{\text{LQG}} = \bigoplus_{\Gamma} \mathcal{H}_{\Gamma}$$

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## References:

- 2004 The Laplacian of a graph as a density matrix: a basic combinatorial approach to separability of mixed states. Braunstein, Ghosh, Severini. quant-ph/0406165
- 2008 The Von Neumann entropy of networks. Passerini, Severini. 0812.2597

$\Gamma = (N, L)$  undirected simple graph

Adjacency  $[A(\Gamma)]_{n,m} = 1$  if  $\{n, m\} \in L(\Gamma)$  and  $[A(\Gamma)]_{u,v} = 0$  otherwise

Degree  $[\Delta(\Gamma)]_{n,n} := d_n = \#$  links adjacent to the node  $n$

Laplacian  $L(\Gamma) := \Delta(\Gamma) - A(\Gamma)$

Density  $\rho_\Gamma := \frac{L(\Gamma)}{d_\Gamma} = \frac{L(\Gamma)}{\text{Tr}(\Delta(\Gamma))}$  Hermitian, positive semi-definite, trace-1

Entropy  $S(\Gamma) = -\text{Tr}[\rho_\Gamma \log \rho_\Gamma]$

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## Plan of the talk

### 1 Interaction Hamiltonian

single quantum particle on the quantum gravitational field  
- construction and properties -

### 2 comparison with Graph Theory

our Hamiltonian provides (up to a certain approximation)  
the same objects already studied in Graph Theory

### 3 Loop Quantum Thermodynamics

this Hamiltonian is a tool to construct statistical objects in LQG



connection variables

Phase Space  $(q_{ab}(x), \pi^{ab}(x), X^a, P_a)$

Hamiltonian constraint  $C(x) = H_{ADM}(x) + \delta^3(x, X)P_0$   $P_0 = \sqrt{P^2 + m^2} \sim m + \frac{P^2}{2m}$

$$P^2 = q^{ab}(x)P_a P_b$$

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$$H = \int dx \delta^3(x, X) N(x) q^{ab}(X) \frac{P_a P_b}{2m}$$

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$$H = \int dx f_R(x, X) \frac{E^{ai}(x)E^{bj}(X)}{\sqrt{q(X)}} \frac{P_a P_b}{2m}$$

pointlike nature of the particle  $\rightarrow$  regularization by a smearing function

$$f_R(x, X) = \begin{cases} \frac{1}{V_R} = \frac{3}{4\pi R^3} & \text{if } |x - X| \leq R \\ 0 & \text{if } |x - X| \geq R \end{cases}$$

## Quantum States

Spin network states  $|s, x\rangle \equiv |s\rangle \otimes |x\rangle \in \mathcal{H}_{\text{LQG}} \otimes \mathcal{H}_P$ .

$$\mathbb{I}_P = \int dx |x\rangle\langle x|$$
$$\langle x|y\rangle = \delta(x, y)$$

the volume operator vanishes everywhere except at the nodes

$$\sqrt{q(x)} |s\rangle = \sum_{n \in N(s)} \nu_n \delta(x, x_n) |s\rangle$$



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## Quantum States

restriction to the nodes

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$$\langle x | y \rangle = \frac{1}{\sqrt{q(x)}} \delta(x, y)$$

$$\langle s, x_n | s', x_{n'} \rangle = \nu_n^{-1} \delta_{ss'} \delta_{nn'} \quad \mathbb{I} = \sum_s \sum_{n \in N(s)} \nu_n |s, x_n\rangle\langle s, x_n|$$

the volume operator vanishes everywhere except at the nodes

$$\sqrt{q(x)} |s\rangle = \sum_{n \in N(s)} \nu_n \delta(x, x_n) |s\rangle$$

## Quantum States

for later convenience...  $|\underline{x}\rangle := \sqrt[4]{q(x)}|x\rangle$ Spin network states  $|s, \underline{x}\rangle \equiv |s\rangle \otimes |\underline{x}\rangle \subset \mathcal{H}_{\text{IQG}} \otimes \mathcal{H}_P$ .

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## Quantum Operators

$$E^{ai}(x) |s\rangle = \kappa \hbar \sum_{\ell} \int_{\ell} dt \dot{\ell}^a(t) \delta^3(x, \ell(t)) |s, \tau^i\rangle$$

$$(\kappa \hbar)^2 \sum_{\ell, \ell'} \int_{\ell} dt \int_{\ell'} dt' \dot{\ell}^a(t) \dot{\ell}^b(t') \delta^3(x, \ell(t)) \delta^3(x, \ell(t')) j_{\ell}(j_{\ell} + 1) |s\rangle =$$

$$P_a = -i\hbar D_a \quad \text{covariant derivative}$$

$$\langle \mathbf{s}, \psi | H | \mathbf{s}, \phi \rangle = \frac{\kappa^2 \hbar^4}{2m} \sum_{\ell} j_{\ell}(j_{\ell} + 1) \int_{\ell} d\mathbf{s} \int_{\ell} dt \overline{\partial_{\mathbf{s}} \psi(\underline{\ell}(\mathbf{s}))} \partial_t \phi(\underline{\ell}(t)) f_R(\underline{\ell}(\mathbf{s}), \underline{\ell}(t))$$

Planck scale!  $\Delta_{\ell} \psi := \int_{\ell} d\mathbf{s} \partial_{\mathbf{s}} \psi(\underline{\ell}(\mathbf{s})) = \psi(\underline{\ell}_f) - \psi(\underline{\ell}_i)$

particle states

$$|\underline{\ell}\rangle := |\underline{\ell}_f\rangle - |\underline{\ell}_i\rangle$$

$$H = \frac{(8\pi\hbar)^2 \ell_{\text{Pl}}^4}{2m v_R} \sum_{\mathbf{s}, \ell \in \mathbf{s}} j_{\ell}(j_{\ell} + 1) |\mathbf{s}, \underline{\ell}\rangle \langle \mathbf{s}, \underline{\ell}|$$

trace-1  $\text{Tr} H = \frac{\hbar^2 \ell_{\text{Pl}}^4}{2m^*} \sum_{n \in N} v_n \sum_{\ell \in n} j_{\ell}(j_{\ell} + 1)$

$$\rightarrow \rho = \frac{H}{\text{Tr} H} = \frac{\sum_{\ell \in n} j_{\ell}(j_{\ell} + 1)}{\sum_{n \in N} v_n \sum_{\ell \in n} j_{\ell}(j_{\ell} + 1)} |\underline{\ell}\rangle \langle \underline{\ell}|$$



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Approximation!  $j_\ell$  all the same  $v_n$  all the same

### Adjacency

$$\begin{aligned} \langle n | \rho | m \rangle_{ph} &= \frac{1}{d_\Gamma} \sum_{\ell} (\delta_{n, \ell_f} - \delta_{n, \ell_i})(\delta_{m, \ell_f} - \delta_{m, \ell_i}) \\ &= \frac{1}{d_\Gamma} (-1) \text{ if } \{n, m\} \in \mathcal{S} \text{ and } [A(\Gamma)]_{u,v} = 0 \text{ otherwise} \end{aligned}$$

### Degree

$$\langle n | \rho | n \rangle_{ph} = \frac{1}{d_\Gamma} \sum_{\ell} (\delta_{n, \ell_f} - \delta_{n, \ell_i})^2 = \frac{d_n}{d_\Gamma}$$

### Entropy

$$S = -\text{Tr}[\rho \log \rho] \equiv S_{BGS}$$

Approximation!  $j_\ell$  all the same  $v_n$  all the same

### Adjacency

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### Entropy

$$S = -\text{Tr}[\rho \log \rho] \equiv S_{BGS}$$



$$\langle E \rangle = \text{Tr}_{\text{LQG}} [H' \tilde{\rho}] = -d(\ln Z)/d\mu \quad \text{mean energy of the particle on a gravitational field}$$

from a measure of the particle  
 (geometry known)

from a measure of the geometry  
 (particle position known)

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Partition Function

$$Z(\mu) = \sum_s e^{-\mu \bar{d}_s}$$

$$\bar{d}_s(\mu) = \frac{\sum_n^N d_n}{N} = \frac{2\ell}{N}$$

Energy Density

$$\rho_s(\mu) = \frac{1}{Z(\mu)} e^{-\mu \bar{d}_s} = e^{-\frac{2}{N}\ell} \left(1 + e^{-\mu \frac{2}{N}}\right)^{-L}$$

$$\langle d \rangle = \frac{1}{Z(\mu)} \sum_s \bar{d}_s e^{-\mu \bar{d}_s} = -\frac{1}{Z(\mu)} \frac{d}{d\mu} Z(\mu) = \frac{2}{N} L \left(1 + e^{-\mu \frac{2}{N}}\right)^{-1}$$

$$\Delta d = \langle d^2 \rangle - \langle d \rangle^2 = -\frac{4L}{N^2} e^{-\mu \frac{2}{N}} \left(1 + e^{-\mu \frac{2}{N}}\right)^{-2}$$

Entropy

$$S = \mu \langle d \rangle - \ln Z(\mu) = \mu \frac{2}{N} L \left(1 + e^{-\mu \frac{2}{N}}\right)^{-1} - L \ln \left(1 + e^{-\mu \frac{2}{N}}\right)$$



Partition Function  $Z(\mu) = \sum_s e^{-\mu \bar{d}_s} = \left(1 + e^{-\mu \frac{2}{N}}\right)^L$

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## Summary

- Hamiltonian of a non-relativistic particle on a gravitational field described by a spinnetwork;
- we have calculated a partition function, and from this other statistical quantities... is this a first step through a viable thermodynamics of the gravitational field?

“Single particle in quantum gravity and BGS entropy of a spin network”  
by C. Rovelli and FV. Phys.Rev.D81:044038,2010 (arXiv:0905.2983)